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Structural correlations and melting of B-DNA fibres

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Despite numerous attempts, the understanding of the thermal denaturation of DNA is still a challenge due to the lack of structural data at the transition since standard experimental approaches to DNA melting are made in solution and do not provide spatial information. We report a measurement using neutron scattering from oriented DNA fibres to determine the size of the regions that stay in the double-helix conformation as the melting temperature is approached from below. A Bragg peak from the B-form of DNA has been observed as a function of temperature and its width and integrated intensity have been measured. These results, complemented by a differential calorimetry study of the melting of B DNA fibres as well as electrophoresis and optical observation data, are analysed in terms of a one-dimensional mesoscopic model of DNA.

I. INTRODUCTION

The X-Ray diagrams published by Wilkins et al. [1] and Franklin et al. [2] in the same issue of Nature as the famous paper of Watson and Crick [3] describing the structure of DNA revealed the significance of the fibre form of DNA in providing oriented samples necessary for structural studies. These images, showing the cross pattern typical of a helix and two strong spots associated to the stacking of the bases in B-DNA, were however illustrating only one aspect of the molecule, its average static structure. In reality the DNA molecule is a highly dynamical object. Its base pairs fluctuate widely. The lifetime of a closed base pair is only of the order of 10 ms [4]. The local opening of the pairs is important for biological function as it allows the reading of the genetic code. When temperature is raised above the physiological range, thermally induced base-pair openings become more cooperative, leading to the so-called “DNA bubbles” i.e. open regions which may extend over tens of base pairs. At sufficiently high temperature they extend over the full molecule and the two strands separate from each other. For a physicist the thermal denaturing of DNA, also known as DNA “melting”, is a phase transition, which is particularly interesting because it occurs in an essentially one-dimensional system. DNA melting started to attract attention soon after the discovery of the double helix structure [5, 6] and was widely studied, providing insights on the interactions within DNA, the influence of base pair sequence on DNA unwinding, and the effect of the solvent on DNA stability [7]. It recently attracted a renewed interest thanks to High Resolution Melting methods (HRM) [8] whereby precise denaturing profiles provide a new tool for biology laboratories.

Despite numerous attempts, the understanding of this remarkable thermodynamic phase transition is still a theoretical challenge. Statistical physics of DNA thermal denaturing has a long history [7] because it raises the fundamental question of a phase transition in a one-dimensional system, but also for practical applications such as the design of Polymerase Chain Reaction (PCR) probes or the HRM studies for biology. The models for DNA denaturation fall in two classes. First, Ising models treat a base pair as a two-state system, which is either closed or open. This is the case of the prototype Poland-Scheraga model [9]. Those models are appealing for their simplicity and because their parameters have been well calibrated. However their drawback is that they need a large number of phenomenological parameters and, for genomic sequences, the calculation can become heavy due to non-local entropy contributions. The second class of models goes beyond a description in terms of two-state systems by incorporating some elements of the structure. The Peyrard-Bishop-Dauxois (PBD) model [10] is still simple because it represents the status of a base pair by a single real number measuring the stretching of the bond between the bases, but contains nevertheless a minimal structural information relevant for structure factor calculations. In its simplest version this model allows a fast calculation of melting curves of long natural DNAs with only 7 parameters [11]. However the success of different models in describing complex melting profiles [11, 12] shows that the correct fit of those curves is not a sufficient test to validate a theory. Examination of further observables, with a more direct link to the underlying structural details, appears necessary. As for other classes of phase transitions in physics, such as magnetic systems, an important feature that characterises the nature of a transition is the growth of the size of the correlated do-

mains as the transition is approached [13].

Traditional methods to investigate the DNA melting transition cannot provide this kind of spatial information. The standard experimental method is to record the sharp increase of UV absorbance at 260 nm which is associated with the un-stacking of the base pairs, while slowly heating a dilute DNA solution. Other approaches rely on circular dichroism measurements or calorimetric studies that measure the heat absorbed at the transition. Although melting curves, showing the fraction of open base pairs versus temperature, may exhibit a multi-step behaviour related to the sequence, none of the experiments are actually sensitive to spatial information, such as the size of the intact regions of the double helix.

This structural information, which is essential to understand the nature of a phase transition, has been lacking. Neutron scattering can provide this missing piece of information, provided the experiment can be performed on an oriented DNA sample. As shown by the historical studies that revealed the double-helical structure [1, 2] this is possible with fibre diffraction. The methods to produce fibre samples have been refined, by controlling the ionic and water content, it is now possible to make high quality fibres with various configurational structures [14, 15].

Here we report a measurement using neutron scattering from oriented DNA fibres (Sec. II), to determine the size of the regions that stay in the double-helix conformation as the melting temperature is approached from below and we show how it can be analysed in terms of the one-dimensional mesoscopic PBD model of DNA [10] (Secs III and IV). We recently published a brief report of those results [16], which are here presented with further data and discussion.

II. EXPERIMENTS

Due to the regular stacking of the base pairs, the B-form of DNA can be viewed as a one-dimensional diffraction grating. This is reflected in a strong Bragg peak for a longitudinal component of the scattering vector $Q_{\parallel} \approx 1.87 \text{ \AA}^{-1}$, associated with the average distance $a = 3.36 \text{ \AA}$ between consecutive base pairs. The principle of our experiment is simple: by following the evolution of this peak as temperature is raised from room temperature to the denaturation temperature T_c , we can monitor the breaking of this “diffraction grating” into pieces separated by denatured regions, where the base stacking is destroyed. We expect a strong broadening of the diffraction peak as T_c is approached. The width of the peak allows us to determine how the average size of the intact double helical domains evolves when DNA approaches its denaturation temperature, which is critical information for a theoretical analysis of the transition. The interest of this method that focuses on a single, intense, diffraction peak is that we precisely collect the information of in-

terest in a measurement which is only weakly perturbed by sample imperfections. In fibres, the B-form of DNA is semicrystalline [15]. The misalignment of the DNA molecules has been estimated to be less than 5 degrees [17]. Its effect on the projection of the base pair distance on the axis is less than 0.5 %, i.e. negligible with respect to other effects such as the variation of the inter base pair distance as a function of the sequence. As a result, for a cut in reciprocal space along the fibre axis, the width of the Bragg peak is not affected. Moreover, by performing a scan off-centre, i.e. with a scattering vector \mathbf{Q} which has a non-zero component orthogonal to the molecular axis, we can also probe the displacement of the base pairs in the transverse direction. Such a scan is not immune from the influence of the misalignment of the molecules which broaden the peak, but it provides interesting data on the fluctuations due to the opening of the base pairs in the vicinity of the thermal denaturation.

A. Materials and methods

Sample preparation

The samples were made from natural DNA extracted from salmon testes (Fluka). The DNA had been oriented using a “spinning” technique [14], whereby the DNA is precipitated out of a 0.4 M lithium salt solution, drawn to a fibre and then wound around a bobbin to make a film of parallel fibres. The samples are then dried, cut from the bobbin, and then stored for a number of weeks in an atmosphere humidified to 75% using $^2\text{H}_2\text{O}$. This fixed the water content of the DNA, ensuring a B-form configurational structure and significantly reduced incoherent neutron scattering from protonated hydrogen in the sample. The B-form was confirmed using x-ray and neutron diffraction.

The neutron scattering samples were folded in concertina fashion, thus preserving the fibre axis direction. The samples were then placed in a niobium envelope and sealed between aluminum plates using lead wire for the seal. The sample cassette was therefore airtight which maintained the water content throughout the experiment. The sample mass was $\sim 1 \text{ g}$.

Neutron scattering

Preliminary diffraction measurements to establish the configurational form were carried out using the IN3 three axis spectrometer at the Institut Laue-Langevin, France. The instrument was configured with a pyrolytic graphite (PG) monochromator and analyser, and the wavelength was set to 2.36 \AA ($\equiv 14.7 \text{ meV}$). The Q-resolution was defined by $60'$ collimation before and after the sample, and higher order wavelength contamination was suppressed using a PG filter. The instrument was used to measure reciprocal space maps, and an example is shown in figure 1 .

The neutron three axis spectrometer IN8, also at ILL, was used to measure the main Bragg peak as a function of temperature. This instrument was configured with

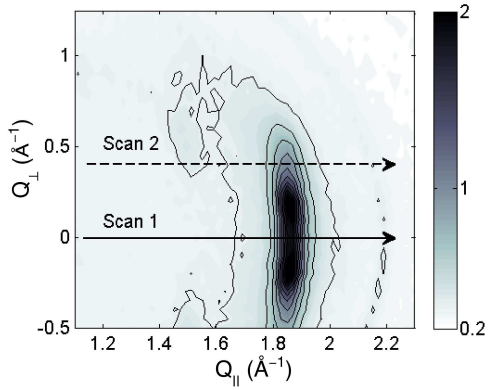


FIG. 1. Reciprocal space map of B-form Li-DNA. The axes are the momentum transfer parallel (Q_{\parallel}) and perpendicular (Q_{\perp}) to the fibre axis. The strong Bragg peak is observed at $Q_{\parallel} = 2\pi/a$ where $a \approx 3.36 \text{ \AA}$ is the average distance between the base pairs along the fibre axis. A powder diffraction peak, coming from the lead wire used to seal the cassette, is just visible at $Q = 2.2 \text{ \AA}^{-1}$. Also marked on the figure are the two standard scans that were repeated at all temperatures. The figure is identical to Fig. 2 [16].

a PG monochromator delivering an incident wavelength of 1.53 \AA ($\equiv 35 \text{ meV}$). The Q -resolution was defined with $40'$ collimation before and after the sample, and was measured by making a reciprocal space map of the (111) Bragg peak from a silicon single crystal. No energy analysis was used, and the static approximation was assumed to hold for the measurements. Temperature control was achieved using a liquid helium cryofurnace.

Two scans were repeated at all temperatures. Their trajectories were calculated for nominally elastic scattering and are shown in figure 1. Scan 1 was along the fibre axis, through the centre of the Bragg peak. The Q_{\perp} for scan 2 was chosen such that, when $Q_{\parallel} = 2\pi/a$, the direction of the scattered beam would be perpendicular to the fibre axis. This type of scan has been used to measure critical phase transitions in low dimensional magnets and assists the static approximation [13]. The temperature steps close to the melting transition were very small (0.1 K) and measurements at a given temperature were repeated numerous times to ensure thermal equilibrium and reproducibility. Examples of the scans at different temperatures are shown in figure 2.

The data were fitted with the Lorentzian function

$$S(Q_{\parallel}) = \int_{-\infty}^{\infty} S(Q_{\parallel}, \omega) d\omega = \frac{I_0}{\pi} \frac{\Gamma/2}{(\Gamma/2)^2 + (Q_{\parallel} - Q_0)^2}, \quad (1)$$

where Q_0 is the peak centre, I_0 is the integrated intensity and Γ is the width. The function was convoluted with the instrument resolution. A second Lorentzian centred at $Q_{\parallel} \approx 1.5 \text{ \AA}^{-1}$ was needed to fit the scan 2 data. The amplitudes for these two peaks were free parameters, however their widths were set to be equal in the fits. Examples of the fits are also shown in figure 2.

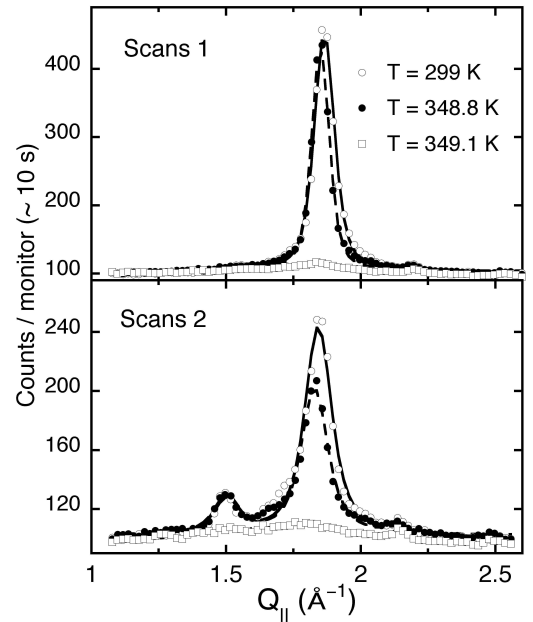


FIG. 2. Examples of the scans shown in figure 1. The data at 299 K represent the starting point for the experiment. The sample is in the melting transition at 348.8 K. The fibre structure has collapsed at 349.1 K. The data have been fitted with equation 1 and the fits are also shown. The figure is identical to Fig. 2 [16].

Calorimetry

Differential Scanning Calorimetry studies have been performed with a DNA film identical to the one used for neutron scattering, but prepared from a different DNA solution. Two different samples were used, with masses of 49 mg and 45.5 mg. The samples cut in the film were rolled and hermetically sealed into the hastelloy sample tube of a Setaram Micro DSC III calorimeter. The reference tube of this differential calorimeter was empty. After a cooling to 278 K the temperature T has been raised to 368 K or 383 K at a rate of 0.6 K/min, maintained for 10 min at the maximum temperature and decreased to 278 K at the same rate of 0.6 K/min. The differential heat flux $\Delta\phi$ has been measured as a function of time (temperature) and the specific heat has been obtained from $(\Delta\phi + \tau \frac{d\Delta\phi}{dt}) / (dT/dt)$ where τ is the thermal time constant of the calorimeter (here $\tau = 60\text{s}$) [19].

Optical observations

A small piece of DNA film has been sealed between two glass plates to preserve its water content while it was heated on a hot plate below an optical microscope at the rate of 1 K/min. The sample was lighted and observed from the top as it went through the denaturation temperature of DNA.

Gel electrophoresis

A small piece of the sample (0.01 g taken before and after heating in the neutron scattering experiment) was

dissolved in water. The solution has been used to run a standard gel electrophoresis experiment, using a 1% non-denaturing agarose gel and stained with Ethidium Bromide. Comparisons with DNA mass ladders were used to measure the length of the DNA fragments in the sample.

A similar experiment was performed with the solution used to prepare the DNA fibres to probe the state of the DNA molecules prior to any treatment.

B. Experimental results

Figures 2 and 3 illustrate the main result given by neutron diffraction. The integrated intensity of the Bragg peak stays constant from room temperature to about $T = 346$ K. At this temperature it starts to show a small decrease occurring on a temperature range of about 3 K, followed by an abrupt drop. A more careful examination of Fig. 3 exhibits the following results:

- The intensity of the peak observed in scan 1 is larger than that of the off-centre peak in scan 2, as expected, but its evolution versus temperature is remarkably similar in both scans.
- The width of the off-centre peak (scan 2) is significantly larger than the width at the centre of the diffraction spot.
- In the 300 – 340 K temperature range the width of both peaks is essentially constant and even shows a slight decrease which can be attributed to an annealing of the sample as shown by Fig. 4. In the experiment shown in this figure, another sample was heated up to 340 K then cooled down to 330 K and heated up to 340 K again. During the first heating the width of the peak decreases, but then, in the cooling stage it keeps its lowest value, as well as during reheating. In this experiment the integrated intensity of the peak stays constant for all temperatures.
- For scan 1, Fig. 3-c shows that, in the vicinity of the transition, where the annealing has been completed, the width of the Bragg peak is remarkably constant until the temperature where the intensity drops abruptly. This width does not show any precursor effect, even when the intensity of the mode starts to decrease in the temperature range $346 < T < 349$ K. On the contrary for scan 2 (off-centre) the width of the peak shows a gradual increase in this temperature range, which appears to mirror the decrease of the intensity of the peak.

The drop of the Bragg-peak intensity, shown in Fig. 3 is extremely sharp in terms of temperature, but requires some time to complete. This is shown in Fig. 5. When the transition was reached and the intensity started to drop, we stopped cooling and observed the time evolution of the peak at constant temperature $T = 349.05 \pm 0.03$ K.

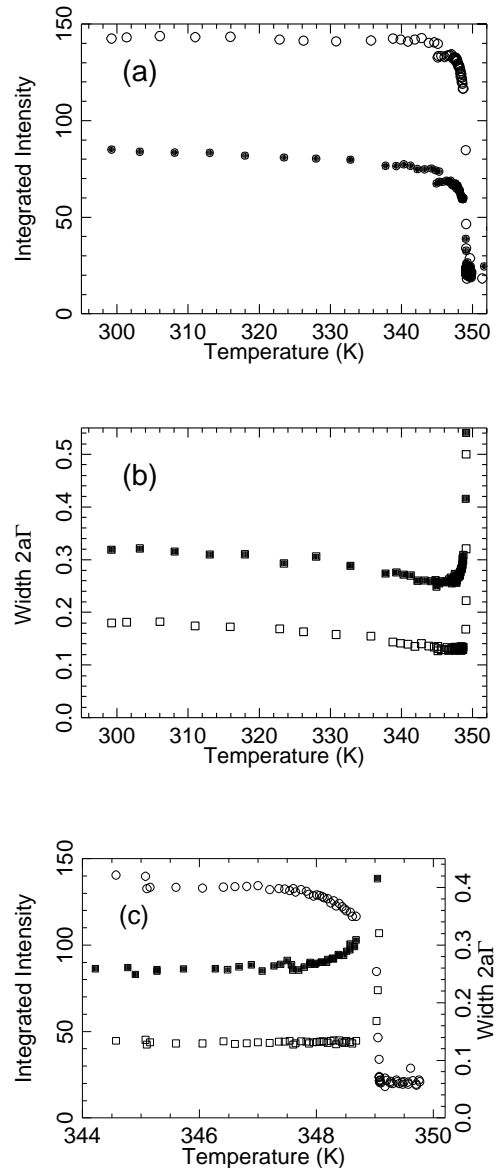


FIG. 3. Integrated intensity (a) and width $2a\Gamma$ (dimensionless) (b) of the Bragg peaks versus temperature. The small discontinuity at $T \sim 345$ K is due to a small misalignment of the instrument that was discovered and subsequently corrected. Results for scan 1 use empty symbols while results for scan 2 are plotted with filled symbols. The bottom panel (c) shows the temperature evolution of the intensity (circles) and width of the peak in the immediate vicinity of the melting transition. The evolution of the intensities are the same for both scans, hence data is only shown for scan 1. The widths are shown for scan 1 (open squares) and scan 2 (closed squares).

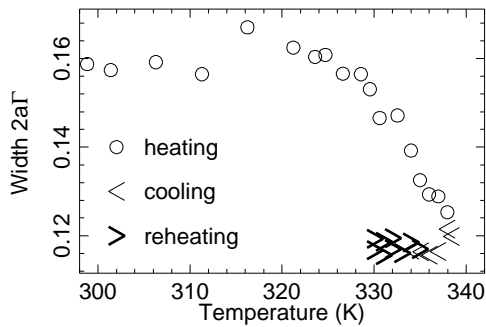


FIG. 4. Experiment showing the annealing of the sample due to heating to moderate temperature. The data correspond to a scan of type 1. The Bragg peak gets sharper on heating (open circles) and keeps its lower width if it is subsequently cooled and heated again.

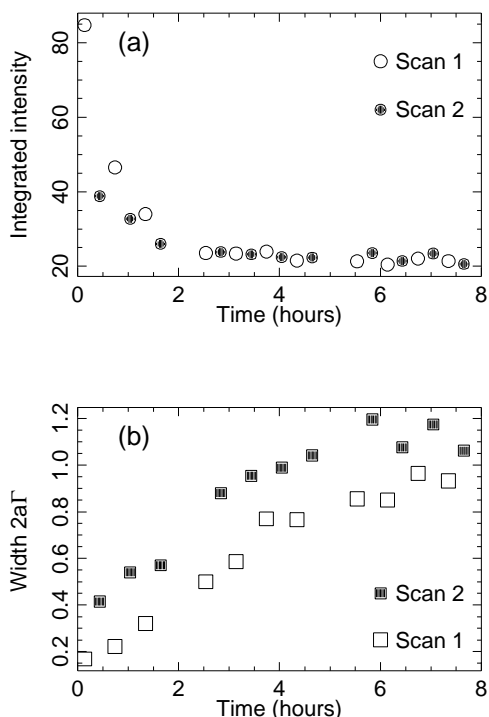


FIG. 5. Time dependence of the integrated intensity (a) and width (b) of the Bragg peak at fixed temperature $T = 349.05 \pm 0.03$ K where the sharp drop of the peak intensity occurs.

Figure 5 shows that the intensity needed about 3 hours to stabilise to a low value, and the width was not stabilised before about 6 hours.

In the fibre sample the transition that corresponds to the almost complete vanishing of the Bragg peak associated to the stacking of the base pairs is not reversible. On cooling we did not observe a reappearance of the peak, although a very small recovery can be observed on a mag-

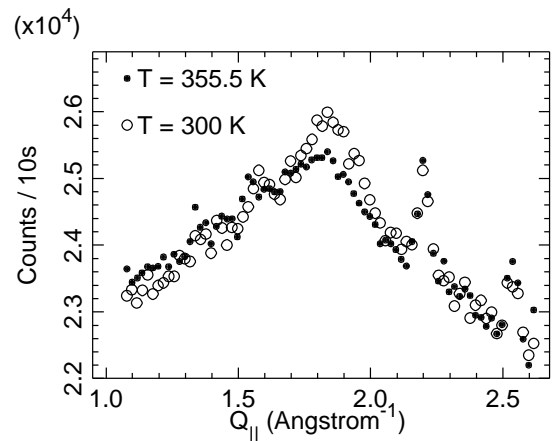


FIG. 6. Plot of the Bragg peak remaining after heating the sample to 355 K (closed symbols) and after subsequent cooling to room temperature (open symbols). The sharp peak around 2.2 \AA^{-1} is due to the aluminium of the sample holder.

nified picture of the Bragg peak recorded at room temperature after cooling as shown in Fig. 6.

Figure 7 shows the specific heat of a DNA film obtained by Differential Scanning Calorimetry (DSC). The sharp peak at 360 K can be attributed to the thermal denaturation of DNA although the transition temperature T_c cannot be quantitatively compared with neutron observations because the measurement was done on a different sample. T_c strongly depends on external conditions, and particularly the ionicity of the solvent [18] so that the shape of the melting curve is more significant than the value of the denaturation temperature when comparing samples. The transition is not reversible, and the sharp peak does not reappear if the sample is cooled and the measurement repeated [20].

The fibre structure of the film is clearly visible in optical microscopy observations of heated DNA films until the denaturation temperature is reached. Then this organised structure of parallel fibres is essentially lost and the film tends to shrink (Fig. 8).

Gel electrophoresis shows that the length of the DNA molecules in the solution used to prepare the film, or in a piece of film which has not been heated, is of the order of 20 kilo-base or larger. However the same measurements performed on a piece of film which has been heated up to DNA denaturation and film collapse, and subsequently cooled to room temperature, only detects DNA fragments of a few hundreds of bases, indicating that the original molecules have been chopped in the thermal cycle (Fig. 9).

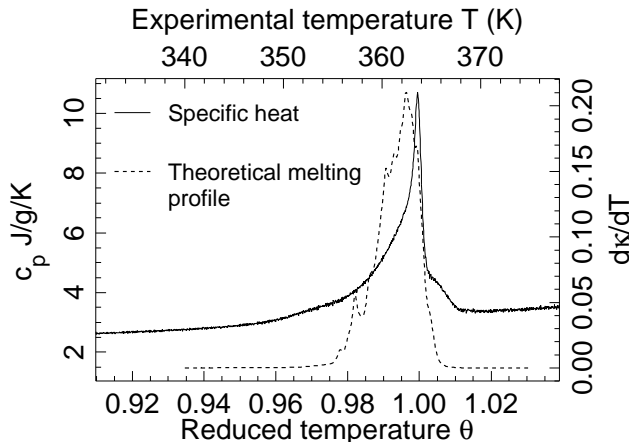


FIG. 7. Specific heat, c_p , of a B-form Li-DNA film similar to the film used in neutron scattering experiments (full line, left scale) obtained by DSC. The dash line shows the theoretical denaturation profile of the sequence used for the analysis (Section III) : derivative with respect to temperature of the fraction κ of open base pairs. The figure is identical to Fig. 1 [16].



FIG. 8. Optical microscopy observation of a piece of film prior (a) and after (b) heating above the denaturing temperature of DNA.

III. ANALYSIS

To analyse the neutron diffraction results, we need to proceed in two steps. First, we must determine the diffraction pattern of a finite segment of double stranded DNA, taking into account the local inhomogeneities in its structure which are associated to base pair sequence, and the thermal fluctuations. Second we must study the statistical physics of DNA to determine the size distribution of the closed segments of DNA as a function of temperature.

A. Structure factor of a closed DNA segment.

We consider a structurally disordered linear chain of M sites. Let a , the average base-pair spacing, be the average distance between successive sites and λ_j the local structural deviation from that value between j th and

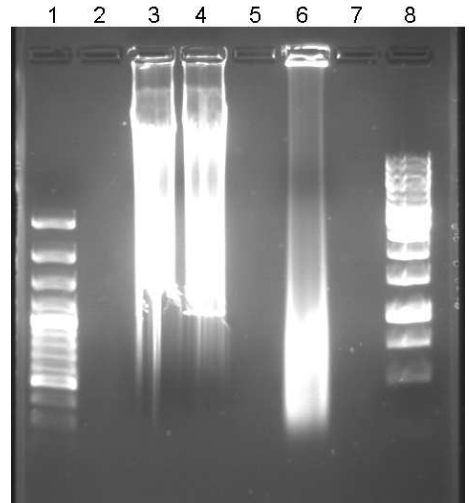


FIG. 9. Electrophoresis image showing the length of the DNA molecules before and after heating of the film. From left to right: DNA mass ladder SM0321 (100-3000 base pairs) (lane 1), solution used to prepare the DNA film (lane 3), solution prepared from a piece of film that has not been heated (lane 4), solution prepared from a piece of film after heating (lane 6), DNA mass ladder SM0311 (250-10000 base pairs) (lane 8). Lanes 2,5,7 were not used.

$(j+1)$ st sites; the equilibrium positions of m th and n th sites differ by $(m-n)a + \sum_{j=n+1}^m \lambda_j$. Structural disorder in the transverse direction is similarly expressed by η_j , the local deviation (from zero) of the distance between j th and $(j+1)$ st sites in the transverse direction. The structure factor of such a finite chain segment is given by

$$S_M(\mathbf{Q}) = \frac{1}{M} \sum_{m,n=1}^M e^{iQ_{\parallel}a(m-n)} \langle e^{i \sum_{j=m+1}^n (Q_{\parallel} \lambda_j + Q_{\perp} \eta_j)} \rangle \times \langle e^{iQ_{\parallel}(x_m - x_n) + iQ_{\perp}(y_m - y_n)} \rangle \quad (2)$$

where x_m, y_m represent, respectively, the longitudinal and transverse displacements of the m th site from its position at thermal equilibrium, and \mathbf{Q} the scattering vector, having the component Q_{\parallel} along the helix axis and Q_{\perp} orthogonal to it. Equation (2) is a slight generalisation of the finite para-crystal theory [21] to account for disorder and motion in the transverse direction. The angular brackets denote averages which can be decoupled because the first refers to structural disorder and the second to thermal motion. To a first approximation, structural disorder is modelled by Gaussian variables $\{\lambda_j\}$ and $\{\eta_j\}$ with zero average and

$$\begin{aligned} \langle \lambda_i \lambda_j \rangle &= \langle \lambda^2 \rangle \delta_{ij} \\ \langle \eta_i \eta_j \rangle &= \langle \eta^2 \rangle \delta_{ij} \\ \langle \eta_i \lambda_j \rangle &= \chi [\langle \eta^2 \rangle \langle \lambda^2 \rangle]^{1/2} \delta_{ij} \end{aligned} \quad (3)$$

where δ_{ij} is the Kronecker symbol ($\delta_{ij} = 1$ if $i = j$ and 0 otherwise). We use estimates of the variances $\langle \lambda^2 \rangle$ and $\langle \eta^2 \rangle$ obtained from conformational analysis [22] and

present alternative calculations for uncorrelated ($\chi = 0$) and correlated ($\chi \neq 0$) structural disorder. Thermal fluctuations in the longitudinal displacements can be calculated in the harmonic approximation; thermal fluctuations in the transverse direction can also be calculated (cf. next subsection for a particular model). We will describe them in an approximate fashion which takes into account the thermal effects due to sequence heterogeneity, i.e. local variations in $\xi_j = \langle y_{j+1} \rangle - \langle y_j \rangle$; then

$$\langle e^{iQ_\perp(y_m - y_n)} \rangle \approx e^{-Q_\perp^2 \langle \xi^2 \rangle |m-n|/2} \quad (4)$$

where $\langle \xi^2 \rangle$ is the variance of the $\{\xi_j\}$'s.

Putting the various terms together and performing one of the two summations [21] results in

$$S_M(\mathbf{Q}) = M + 2 \sum_{n=1}^{M-1} (M-n) \cos(Q_\parallel n a) e^{-n\Delta}, \quad (5)$$

where

$$\begin{aligned} 2\Delta = & Q_\parallel^2 (\langle \lambda^2 \rangle + \langle \sigma^2 \rangle) + Q_\perp^2 (\langle \eta^2 \rangle + \langle \xi^2 \rangle) \\ & + 2Q_\parallel Q_\perp \chi (\langle \eta^2 \rangle \langle \lambda^2 \rangle)^{1/2} \end{aligned} \quad (6)$$

and $\sigma^2 = k_B T / \mu c_0^2 a^2$, the Debye-Waller correction due to longitudinal thermal motion at any temperature T , can be obtained from the total DNA mass per base pair $\mu = 618$ a.m.u. and the measured [23] sound velocity $c_0 = 2830$ m.s⁻¹; k_B is the Boltzmann constant.

Near the first Bragg peak, which is where the present experiment focused, and for sufficiently large cluster sizes, the sum (5) can be approximated by

$$S_M^*(\mathbf{Q}) \sim M S'(\mathbf{Q}) \doteq M \frac{\sinh \Delta}{\cosh \Delta - \cos(Q_\parallel a)}, \quad (7)$$

B. Statistical physics of the closed regions of DNA.

To calculate the size distribution of the closed segments of DNA by a statistical physics analysis, we selected the PBD model [10] which is sufficiently simple to allow the analysis of DNA segments of thousands of base pairs, but nevertheless includes some data on the structure of the molecule which are necessary to calculate the structure factor. Moreover, as it describes the molecule in terms of a Hamiltonian, its parameters are directly linked to physical quantities. Therefore they are easier to determine than for Ising models, although they still need to be refined by comparison with a variety of experimental melting profiles.

The configuration energy of a DNA molecule of N base pairs is written as

$$H_y = \sum_{j=1}^{N-1} W(y_j, y_{j+1}) + \sum_{j=1}^N V_j(y_j) \quad (8)$$

where y_j represents the stretching of the j^{th} base pair, due to the transverse displacements of the bases. The stacking interaction between adjacent bases is described by the anharmonic potential [10]

$$W(y_j, y_{j+1}) = \frac{1}{2} k \left[1 + \rho e^{-b(y_j + y_{j+1})} \right] (y_j - y_{j+1})^2 \quad (9)$$

which takes into account the weakening of the interactions when the pairs are broken. The potential $V_j(y_j) = D_j [1 - \exp(-\alpha_j y_j)]$ is a Morse potential which describes the combined effects of hydrogen-bonding, electrostatic interactions between the charged phosphate groups, and solvent effects on the the j^{th} base pair. The 4 possible bases, A, T, G, C form two types of pairs, $A - T$, linked by two hydrogen bonds, and $G - C$, linked by three hydrogen bonds. Both the stacking interactions W and the intra-pair potential V are affected by the sequence of bases. However, although subtle sequence effects on short DNA fragments may require the introduction of different stacking potentials for different interacting pairs [24], the melting curves of long DNA molecules, with several thousands base pairs, can be accurately reproduced by introducing the effect of the sequence in the intra-pair potential V_j only [11], which drastically reduces the number of model parameters. Therefore in our calculations the stacking interaction is treated as homogeneous.

To determine the melting curve or calculate the size of the closed regions, we need to give a quantitative definition of a closed base pair. This can be done by choosing a reference stretching y_c . Base pair j is considered as closed if $y_j < y_c$. We select $y_c = 1.5$ Å, which corresponds to a base pair whose stretching is well on the plateau of the Morse potential. The results are weakly dependent of the value of y_c provided it is larger than 1 Å because molecular dynamics simulations show that, once a base pair as been stretched to a value that brings it on the plateau of the Morse potential, it is likely to open widely.

The statistical weight of a given configuration of the molecule thermalised at temperature T is

$$\mathcal{Z} = \int_{y_{l_1}}^{y_{m_1}} dy_1 \int_{y_{l_2}}^{y_{m_2}} dy_2 \dots \int_{y_{l_N}}^{y_{m_N}} dy_N e^{-H_y(y_1, y_2, \dots, y_N)/(k_B T)}. \quad (10)$$

The limits of the lower and upper bounds of the integrals y_{l_j}, y_{m_j} depend on the particular configuration. Setting $y_{l_j} = -\infty, y_{m_j} = +\infty$ for all j allows the molecule to explore its full configurational space and \mathcal{Z} is then the partition function Z . Setting $y_{l_j} = -\infty, y_{m_j} = y_c$ defines a configuration in which base pair j is closed while $y_{l_j} = y_c, y_{m_j} = +\infty$ defines a configuration where it is open. The integrals associated to all those configurations can be easily calculated because the model is one-dimensional and restricted to nearest-neighbour coupling so that, instead of a highly multidimensional integral, one has to compute a chain of one-dimensional integrals [11, 27, 28] involving a kernel which depends on the stretching at two adjacent

sites. Moreover the speed of the calculation can be significantly increased by expanding the site-dependent kernel on the basis of the eigenfunctions of a reference kernel, for instance the kernel associated to an $A - T$ base pair [11].

These calculations allow us to obtain the probability $\mathcal{P}(m, i)$ that m adjacent sites, starting at site i are closed,

$$\mathcal{P}(m, i) = \frac{1}{\mathcal{Z}} \mathcal{Z}(y_1, y_2, \dots, y_{i-1}, y_i < y_c, y_{i+1} < y_c, \dots, \dots, y_{i+m-1} < y_c, y_{i+m}, \dots), \quad (11)$$

by computing the statistical weight of configurations where restrictions on the integration range are imposed for all sites belonging to that closed region and no restrictions are imposed elsewhere. Then these quantities give the probability $P(m, i)$ that a closed cluster of size m , with open ends, exists at site i through

$$P(m, i) = \mathcal{P}(m, i) - \mathcal{P}(m + 1, i - 1) - \mathcal{P}(m + 1, i) + \mathcal{P}(m + 2, i - 1) \quad (12)$$

from which the probability to have a closed cluster of size m in a DNA segment of N_0 base pairs is simply

$$P(m) = \frac{1}{N_0} \sum_{i=1}^{N_0} P(m, i). \quad (13)$$

The average size of a cluster of closed base pairs is then

$$l_c = \frac{\sum m P(m)}{\sum P(m)} = \frac{h}{1 - h - P_0}, \quad (14)$$

where h is the helix fraction and P_0 the statistical weight of two consecutive bases being in the closed state. Similar calculations for the open regions of the DNA molecule give the average size of the denatured regions $l_o = (1 - h)/(1 - h - P_0)$. These quantities are computed for a range of cluster sizes from 1 to a maximum value M . To avoid end effects we study a DNA segment of size $N = N_0 + 2M$ and the sites i are chosen so that the clusters that we consider are formed of the bases pairs which are at least M sites away from the ends.

For a natural DNA sample, which may contain millions of base pairs, the intensity observed in a neutron scattering experiment is proportional to

$$S(\mathbf{Q}) = \sum_{m=1}^{\infty} P(m) S_m(\mathbf{Q}) \quad (15)$$

where $S_m(\mathbf{Q})$ is given by Eq. (5), and the size of the DNA molecule has been extrapolated to infinity.

In practice the calculation is performed with segments of natural DNA which include 280000 base pairs, and $P(m, i)$ is computed up to cluster sizes M of a few hundreds of base pairs (typically $M = 150$). It is however

easy to determine $P(m)$ for large m because it scales exponentially with m [29], as shown in Fig. 10-b. Fitting the numerical data of $\log(P(m))$ for $40 \leq m \leq M$ by a straight line we get $P(m) = P_0 \zeta^m$ for $m > M$ so that in practice the calculation of $S(\mathbf{Q})$ by Eq. (15) is expressed as

$$S(\mathbf{Q}) = \sum_{m=1}^M P(m) S_m(\mathbf{Q}) + \sum_{m=M}^{M_0} P_0 \zeta^m S_m(\mathbf{Q}) + \mathcal{S}'(\mathbf{Q}) P_0 \zeta^{M_0+1} \left[\frac{M_0}{1 - \zeta} + \frac{1}{(1 - \zeta)^2} \right], \quad (16)$$

where the summation for $m > M_0$ ($M_0 = 1000$) has been calculated analytically using the property that, for large m , $S_m(\mathbf{Q})$ can be approximated by the limiting form (7). The same method can be used to compute the average cluster size l_c (Eq. (14)) versus temperature. A typical result is shown on Fig. 10-c. Then the structure factor is fitted with the same Lorentzian expression as the one used to analyse experimental data to determine the integrated intensity I_0 and width Γ of the diffraction peak.

C. Model parameters

To analyse of the neutron scattering experiments, in principle we would need to know the base-pair sequence in the sample. As the experiment requires a significant amount of DNA it can only be performed with natural DNA. The salmon testes DNA that we use is provided without its sequence [30], and even its GC content is only approximately known. It is estimated to be 41.2%.

The theoretical analysis has been tested on different DNA sequences from the genome of *Danio rerio* (zebrafish) [31] and *Pyrococcus abyssi* [32]. The results presented in Fig. 10 have been obtained with a sequence of 280000 bases, part of the full genome of *Pyrococcus abyssi*, chosen because its denaturation curve is the closest to the denaturation profile of our samples measured by differential scanning calorimetry (Fig. 7). The GC content of this fragment is 44.08%

Model parameters have been obtained from an extensive study of DNA denaturation on various sequences [11], which determined a set of parameters allowing the prediction of melting curves of various DNA sequences to a good accuracy. Experiments to record DNA denaturation curves are generally performed in solution with sodium salt, the Na^+ ions being necessary to stabilise DNA by screened the repulsions between the charged phosphate groups. However our neutron scattering experiments have been performed with Li-DNA because its secondary structure is more stable, and stays in the B-form over a much higher humidity range, than that of Na-DNA [25, 26]. Therefore the parameters of the AT and GC Morse potentials cannot be determined unambiguously. The values that we use correspond to a high

sodium content. They have been chosen on the basis of the shape of the melting curve rather than the value of the denaturation temperature T_c . This is why the comparison between theory and experiments must be made with the reduced temperature $\theta = T/T_c$ rather than with the actual temperature. The parameters selected for the analysis are $k = 0.00045 \text{ eV/\AA}^2$, $\rho = 50$ and $b = 0.20 \text{ \AA}^{-1}$, $\alpha_{AT} = 4.2 \text{ \AA}^{-1}$, $\alpha_{GC} = 6.9 \text{ \AA}^{-1}$ and $D_{AT} = 0.12905 \text{ eV}$, $D_{GC} = 0.16805 \text{ eV}$.

D. Results

Figure 10 summarises the main results of the theoretical analysis by showing the variation versus temperature of the integrated intensities and widths of the Bragg peaks in scan 1 and 2, as they are predicted by the model. As expected the intensity shows a sharp drop near the denaturation transition. It reflects the openings of the base pairs that break the clusters of stacked pairs giving rise to the investigated Bragg peak and therefore reduces the number of scattering sites. As a result the integrated intensity of the peak almost provides a quantitative measure of the helix fraction of DNA because, as shown by Eq. (7), for sufficiently large clusters, the structure factor is proportional to the number of sites in a cluster. Moreover Fig. 10-c shows that the size of the clusters significantly drops only in the last stage of the denaturation.

The width of the Bragg peak provides the spatial information that standard observations of DNA denaturation cannot give. It is strongly sensitive to the distribution of the sizes of the diffracting clusters. The drop of the average cluster size in the last stage of the transition (Fig. 10-c) is reflected in the large increase of the width of the Bragg peak predicted by the theory in the high temperature range (Fig. 10-d). For scan 2, with a nonzero transverse component Q_{\perp} of the scattering vector, the width of the peak is also affected by the transverse structural disorder due to the effect of the sequence (variables η_j) and by their correlations with the longitudinal structural disorder (variables λ_j), measured by the coefficient χ in Eq. (3). The statistical properties of λ_j and η_j have been obtained by conformational analysis [22] but their correlations have not been determined. We show results with $\chi = 0$ (no correlation) and $\chi = 0.35$ corresponding to moderate correlations. Moreover scan 2, is also probing the transverse fluctuations of the bases prior to opening. The inset in Fig. 10-d shows that, in the vicinity of the transition, these fluctuations are expected to cause an extra increase of the width of the Bragg peak, for scan 2 only. However this effect is small because the Bragg peak is only generated by the closed sections of the DNA molecules, where the fluctuations are therefore limited to small amplitude motions of the bases.

IV. DISCUSSION

To allow a quantitative comparison of the experimental and theoretical results we plot the data as a function of a reduced temperature $\theta = T/T_c$, where T_c is the temperature where 50% of the bases are open. This is done in Fig. 11. To eliminate the experimental factor associated to the apparatus, in this figure the experimental intensity of the Bragg peak has been rescaled to 1 in the low θ limit, which is also the limit of the theoretical intensity in the low temperature range. The widths have been multiplied by the mean separation of the base pairs, a , to create dimensionless variables. Therefore a quantitative comparison between theoretical and experimental widths is possible and it does not involve any arbitrary factor.

Let us first examine the integrated intensity of the Bragg peak. Both the experimental data and the theoretical curve show an intensity which stays almost constant up to temperatures very close T_c , $\theta \approx 0.97$. As previously noted this reflects the very low fraction of open base pairs until the vicinity of the transition is reached, in agreement with the theoretical result of Fig. 10-a. In the early stage of the transition the theoretical curve follows the experimental decay of the intensity of the peak, but in the immediate vicinity of T_c ($\theta \approx 1$) the experimental intensity shows an almost discontinuous drop, while the theoretical curve has a narrow but smooth decay. This discrepancy is a sign that another phenomenon, not included in the theoretical description happens. As discussed in Sec. II the optical observations of a sample during heating indicate that, at high temperature, the film itself shows an irreversible ‘‘collapse’’ characterised by a disorganisation of the oriented fibre structure. Although this is not described by the model, the theoretical analysis shines some light on this phenomenon. Figure 10-c shows that, above T_c , the length of the denatured regions grows very quickly with temperature, showing almost a divergence. As the single strands are very flexible they can gain a lot of entropy by fully losing their initial orientation so that the remaining double-helix segments are embedded in a liquid-like medium of entangled single strands which quickly becomes the dominant phase in the sample. The rapid growth of the size of the open fragments gives a lot of freedom to the rigid double-helix segments allowing them to lose their spatial orientation, which causes the sharp drop of the intensity of the Bragg peak. At higher temperatures the sample is no longer a good approximation of a one-dimensional crystal, but is more an ensemble of disoriented, short length, DNA. This temperature range corresponds to the shaded area in Figs. 10 and 11. It is interesting to notice that, when this ‘‘collapse’’ of the film occurs, the theory predicts that the size of the closed segments is still large, of the order of 80 – 100 base pairs. This is indirectly confirmed by the electrophoresis analysis of the length of the DNA fragments in the film before and after heating which indicate that, after film melting, the DNA molecules that were

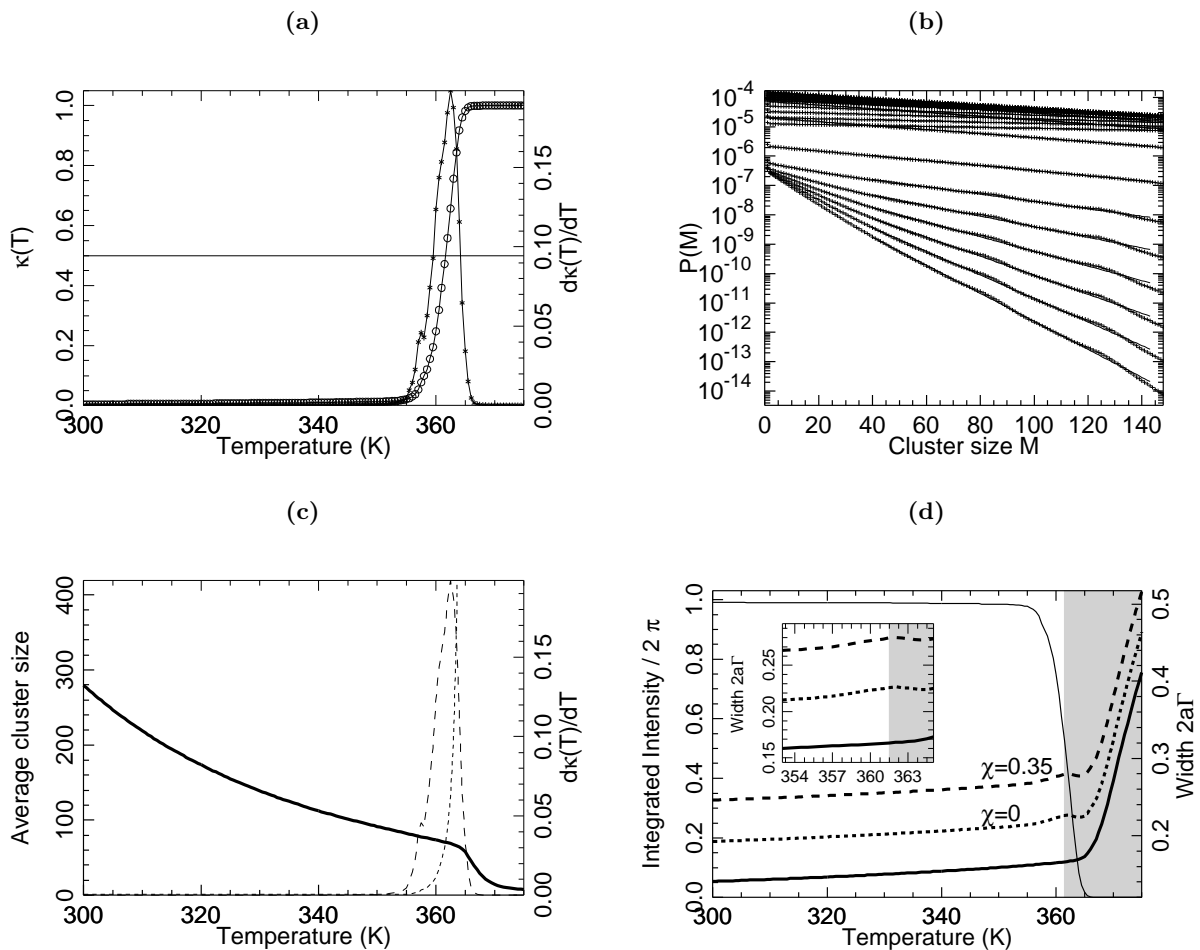


FIG. 10. Theoretical results. (a) Melting curve of a reference DNA segment of 280000 base pairs, part of the genome of *Pyrococcus abyssi*: the circles show the open fraction versus temperature and the stars correspond to its derivative to get the melting profile. (b) Probabilities $P(m)$ versus m at different temperatures in the range $T = 330$ K to $T = 375$ K, in logarithmic scale. The points are the values calculated from the statistical mechanics of the DNA sample, and the lines show a linear fit for $40 \leq m \leq 150$. Curves are plotted every 10 K from 300 to 340 K and every 1.5 K above. (c) The average size the closed clusters l_c (full line) and the average size of the open regions l_o (dotted line) versus T . The dashed line shows the melting profile on the same temperature scale. (d) Integrated intensity (thin full line) and width (thick full line for scan 1, dotted and dashed lines for scan 2) versus temperature. For scan 2 the figure shows two cases: $\chi = 0$, i.e. ignoring the correlation between longitudinal λ_j and transverse η_j components of the structural disorder due to the sequence (dotted line) and $\chi = 0.35$ (dashed line) which assumes a partial correlation. The shaded area of the plot shows the temperature range in which the experimental observations are hindered by the melting of the sample film (Sec. IV). The inset shows a magnification of the variation of the widths versus temperature in the immediate vicinity of the transition.

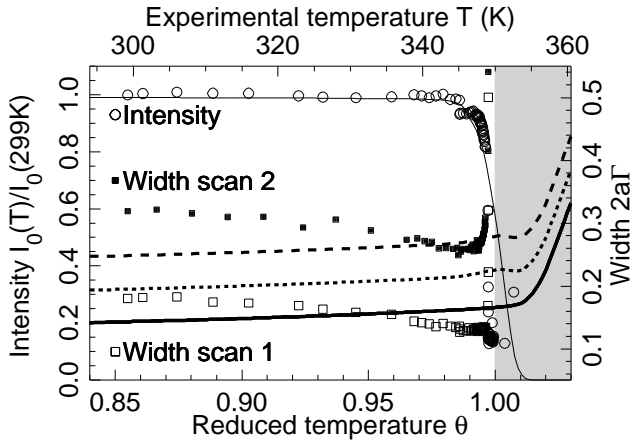


FIG. 11. Comparison between theory and experiment. The points are the experimental results while the curves plot the theoretical results. A reduced temperature $\theta = T/T_c$ is used, where T_c is temperature where 50% of the bases are open. The circles show the integrated intensity of the Bragg peak, rescaled to 1 at low temperature. After rescaling, the results for scans 1 and 2 exactly follow the same curve. The thin line is the calculated integrated intensity. The open squares show the experimental width of scan 1 and the thick full line is the theoretical value for this width. The filled squares show the width of scan 2. The dotted line shows the theoretical width of scan 2 calculated by assuming that the longitudinal and transverse structural disorder due to the sequence, determined by conformational analysis [22], are uncorrelated ($\chi = 0$) while the dash line is the theoretical width of scan 2 calculated by assuming a partial correlation between the longitudinal and transverse structural disorder ($\chi = 0.35$). The shaded area of the plot shows the temperature range in which the experimental observations are hindered by the collapse of the sample film. The figure is identical to Fig. 3 [16].

more than 20000 base-pair long at low temperature are chopped into segments of a few hundreds of base pairs. This breaking can be understood by the high stress concentration that occurs at the end of the rigid fragments, linked to each other by the flexible single strands, when they rotate as the film melts. It is therefore not surprising to detect DNA fragments which have a length of the order of the size of the closed clusters.

Let us now examine the width of the Bragg peak. For scan 1 the calculation of the width does not involve any free parameter once the model has been calibrated to match the denaturation curve of DNA. The other parameters entering in the structure factor calculation (Eqs. (5) and (16)) are derived from the structure of DNA [22] and its sound velocity measured along the helix axis [23]. For scan 1 ($Q_{\perp} = 0$), the calculation gives a result which is in good agreement with experiments (Fig. 11), although the experimental results are probably affected by some annealing of the sample causing a slight decay of the width of the peak (Fig. 4) that the theory does not describe. In

spite of this limitation two points emerge from the comparison between theory and experiments. First the structural data which enter in the calculation of the width for scan 1, and particularly the fluctuations of the base-pair distances along the helix, measured by $\langle \lambda^2 \rangle$, corresponding to a standard deviation of 0.18 \AA , are here tested on a large scale since the width of the diffraction peak involves an average over the billions of base pairs of the sample. The discrepancy of less than 15% between the calculated and experimental widths indicates that the results of the conformational analysis [22] are accurate. Second, for scan 1, the width of the peak is remarkably constant until T_c and the collapse of the film. This indicates that the base pair openings, which start to be very significant at $\theta = T/T_c = 0.98$ do not cause a sharp decrease of the size of the diffracting clusters until the denaturation has occurred. This is what the theoretical model indicates (Fig 10-c). Clusters of about 100 base pairs remain intact well within the denaturation region and denature as a whole, this being allowed by the surrounding open bubbles in the last stage of the denaturation. This provides a good test of the statistical physics description of DNA that we use, validating the model beyond its ability to predict melting curves.

Contrary to scan 1 the calculation of the the width of the Bragg peak in scan 2, with a transverse component Q_{\perp} of the scattering vector, involves an unknown parameter, the coefficient χ that measures the correlation between the longitudinal and transverse structural disorder due to the sequence (Eq. 3). Figure 11 shows that, if we ignore this correlation by setting $\chi = 0$, the theoretical value is about 30% lower than the experimental width. Setting $\chi = 0.35$, i.e. a moderate correlation, we get a theoretical width which matches the experimental value for $\theta \approx 0.97$ which we take as a temperature where the sample is “annealed”. The results suggest that neutron scattering could be used to probe this structural property of DNA and it would be interesting to test this result by conformational analysis.

In the range $0.99 < \theta < 1$ the experiment detects a significant rise of the width of scan 2 which is not shown by the theoretical curve. According to Eq. (5), the transverse fluctuations of the base pairs, which become large because they start to open, bring an extra contribution to the width through a growth of $\langle \xi^2 \rangle$. This contribution is visible on the inset of Fig. 10, but, as discussed in Sec. III, this effect has to be small since only fluctuations in the closed clusters of base pairs can contribute to the shape of the Bragg peak. This is not enough to account for the observed increase of the width of scan 2 below the transition. However there is another contribution to the width which is not included in the structure factor calculation, it is the misalignment of the molecules. It is very likely that the collapse of the film is preceded by increased fluctuations in the orientation of the helix fragments. This should have a strong influence on the width of the Bragg peak in scan 2. For instance orientational fluctuations of 10 degrees which change the projection of

the base pair distance on Q_{\parallel} by less than 2% lead to a projection of 17% of this distance along Q_{\perp} . Only a theory of the collective effects leading to the melting of the film could properly account for this effect.

V. CONCLUSION

In conclusion, we have shown that neutron scattering can be used to monitor the thermal denaturation of DNA, providing the spatial information that other methods cannot measure. By focusing the study on the Bragg peak which is associated to the base pair stacking we can obtain accurate results which are not limited by the fibre nature of the samples. The width of the Bragg peak can be described by a simple nonlinear model for DNA at the scale of base pairs, thus providing further validation of this model which had already proved able to predict complex DNA denaturation curves with a small number of parameters. Moreover, by selecting a scatter-

ing vector which is not parallel to the axis of the DNA helix, the shape of the Bragg peak is also sensitive to the transverse fluctuations of the base pairs, that this dynamical model can calculate. This should allow further comparison between theory and experiments by investigating not only the opening of the base pairs, but also their large scale fluctuations, important in many biological processes. This aspect could not be investigated in the present experiments due to the collapse of the fibre structure of the sample.

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