



**HAL**  
open science

# Multi-dimensional scalar conservation laws with unbounded integrable initial data

Denis Serre

► **To cite this version:**

Denis Serre. Multi-dimensional scalar conservation laws with unbounded integrable initial data. 2018.  
ensl-01837034v2

**HAL Id: ensl-01837034**

**<https://ens-lyon.hal.science/ensl-01837034v2>**

Preprint submitted on 26 Jul 2018

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Multi-dimensional scalar conservation laws with unbounded integrable initial data

Denis Serre  
École Normale Supérieure de Lyon\*

## Abstract

We discuss the minimal integrability needed for the initial data, in order that the Cauchy problem for a multi-dimensional conservation law admit an entropy solution. In particular we allow unbounded initial data. We investigate also the decay of the solution as time increases, in relation with the nonlinearity.

The main ingredient is our recent theory of divergence-free positive symmetric tensor. We apply in particular the so-called *compensated integrability* to a tensor which generalizes the one that L. Tartar used in one space dimension. It allows us to establish a Strichartz-like inequality, in a quasilinear context.

This program is carried out in details for a multi-dimensional version of the Burgers equation.

**Notations.** When  $1 \leq p \leq \infty$ , the natural norm in  $L^p(\mathbb{R}^n)$  is denoted  $\|\cdot\|_p$ , and the conjugate exponent of  $p$  is  $p'$ . The total space-time dimension is  $d = 1 + n$  and the coordinates are  $x = (t, y)$ . In the space of test functions,  $\mathcal{D}^+(\mathbb{R}^{1+n})$  is the cone of functions which take non-negative values. The partial derivative with respect to the coordinate  $y_j$  is  $\partial_j$ , while the time derivative is  $\partial_t$ . Various finite positive constants that depend only the dimension, but not upon the solutions of our PDE, are denoted  $c_d$ ; they usually differ from one inequality to another one.

---

\*U.M.P.A., UMR CNRS-ENSL # 5669. 46 allée d'Italie, 69364 Lyon cedex 07. France.  
denis.serre@ens-lyon.fr

# 1 Introduction

Let us consider a scalar conservation law in  $1 + n$  dimensions

$$(1) \quad \partial_t u + \sum_{i=1}^n \partial_i f_i(u) = 0, \quad t > 0, y \in \mathbb{R}^n.$$

We complement this equation with an initial data

$$u(0, y) = u_0(y), \quad y \in \mathbb{R}^n.$$

The flux  $f(s) = (f_1(s), \dots, f_n(s))$  is a smooth vector-valued function of  $s \in \mathbb{R}$ . We recall the terminology that an entropy-entropy flux pair is a couple  $(\eta, q)$  where  $s \mapsto \eta(s)$  is a numerical function,  $s \mapsto q(s)$  a vector-valued function, such that  $q'(s) \equiv \eta'(s)f'(s)$ . The Kruzhkov's entropies and their fluxes form a one-parameter family:

$$\eta_a(s) = |s - a|, \quad q_a(s) = \operatorname{sgn}(u - a)(f(u) - f(a)).$$

Together with the affine functions, they span the cone of convex functions.

We recall that an *entropy solution* is a measurable function  $u \in L^1_{\text{loc}}([0, +\infty) \times \mathbb{R}^n)$  such that  $f(u) \in L^1_{\text{loc}}([0, +\infty) \times \mathbb{R}^n)$ , which satisfies the Cauchy problem in the distributional sense,

$$(2) \quad \int_0^\infty dt \int_{\mathbb{R}^n} (u \partial_t \phi + f(u) \cdot \nabla_y \phi) dy + \int_{\mathbb{R}^n} u_0(y) \phi(0, y) dy = 0, \quad \forall \phi \in \mathcal{D}(\mathbb{R}^{1+n}),$$

together with the entropy inequalities

$$(3) \quad \int_0^\infty dt \int_{\mathbb{R}^n} (\eta_a(u) \partial_t \phi + q_a(u) \cdot \nabla_y \phi) dy + \int_{\mathbb{R}^n} \eta_a(u_0(y)) \phi(0, y) dy \geq 0, \quad \forall \phi \in \mathcal{D}^+(\mathbb{R}^{1+n}), \forall a \in \mathbb{R}.$$

When it enjoys higher integrability, an entropy solution is expected to satisfy additional entropy inequalities of the form

$$(4) \quad \int_0^\infty dt \int_{\mathbb{R}^n} (\eta(u) \partial_t \phi + q(u) \cdot \nabla_y \phi) dy + \int_{\mathbb{R}^n} \eta(u_0(y)) \phi(0, y) dy \geq 0, \quad \forall \phi \in \mathcal{D}^+(\mathbb{R}^{1+n}),$$

for more general convex entropies  $\eta$ . In particular, one is interested in inequality (4) for the entropy-entropy flux pair

$$\bar{\eta}(s) = \frac{1}{2} s^2, \quad \bar{q}(s) = \int_0^s z f'(z) dz.$$

The theory of this Cauchy problem dates back to 1970, when S. Kruzhkov [7] proved that if  $u_0 \in L^\infty(\mathbb{R}^n)$ , then there exists one and only one entropy solution in the class

$$L^\infty(\mathbb{R}_+ \times \mathbb{R}^n) \cap C(\mathbb{R}_+; L^1_{\text{loc}}(\mathbb{R}^n)).$$

The operator  $S_t : u_0 \mapsto u(t, \cdot)$ , which maps  $L^\infty(\mathbb{R}^n)$  into itself, enjoys several additional properties. On the one hand, a comparison principle says that if  $u_0 \leq v_0$ , then  $S_t u_0 \leq S_t v_0$ . For instance, the solution  $u$  associated with the data  $u_0$  is majorized by the solution  $\bar{u}$  associated with the data  $(u_0)_+$ , the positive part of  $u_0$ . On another hand, if  $v_0 - u_0$  is integrable over  $\mathbb{R}^n$ , then  $S_t v_0 - S_t u_0$  is integrable too, and

$$(5) \quad \int_{\mathbb{R}^n} |S_t v_0 - S_t u_0|(y) dy \leq \int_{\mathbb{R}^n} |v_0 - u_0|(y) dy.$$

Finally, if  $u_0$  belongs to some  $L^p(\mathbb{R}^n)$  space, then  $S_t u_0$  has the same integrability, and the map  $t \mapsto \|S_t u_0\|_p$  is non-increasing. We warn the reader that the contraction property (5) occurs only for the  $L^1$ -norm, but not for other  $L^p$ -norms.

Because of (5) and the density of  $L^1 \cap L^\infty(\mathbb{R}^n)$  in  $L^1(\mathbb{R}^n)$ , the family  $(S_t)_{t \geq 0}$  extends in a unique way as a continuous semi-group of contractions over  $L^1(\mathbb{R}^n)$ , still denoted  $(S_t)_{t \geq 0}$ . When  $u_0 \in L^1(\mathbb{R}^n)$  is unbounded, we are thus tempted to declare that  $u(t, y) := (S_t u_0)(y)$  is the *abstract solution* of the Cauchy problem for (1) with initial data  $u_0$ . An alternate construction of  $(S_t)_{t \geq 0}$ , based upon the Generation Theorem for nonlinear semigroups, is due to M. Crandall [2], who pointed out that it is unclear whether  $u$  is an entropy solution, because the local integrability of the flux  $f(u)$  is not guaranteed. It is therefore an important question to identify the widest class of integrable data for which  $u$  is actually an entropy solution of (1).

To achieve this goal, we develop a new strategy, based on the Compensated Integrability that we introduced in our previous papers [10, 11]. It uses a map  $a \mapsto M(a) \in \mathbf{Sym}_d$ , whose lines are entropy-entropy flux pairs, where the entropies are precisely the functions  $\text{id}_{\mathbb{R}}, f_1, \dots, f_n$  which appear in the conservation law. The map  $M$  is a non-decreasing function of  $a$ . This tensor was already used when  $n = 1$  by L. Tartar [13] to prove the compactness of the semi-group, and by F. Golse [5] (see also [6]) to prove some kind of regularity. An essential ingredient is the amount of non-linearity displayed by the flux  $f$ . We illustrate our strategy by carrying out the details on the most typical nonlinear conservation law, a multi-d generalization of the Burgers equation.

**Outline of the article.** We begin with a detailed, definitive, analysis of the multi-d Burgers equation. The equation is described in the next section. Our main result is a well-posedness

when the initial data is integrable. It is based on a dispersion estimate, which has the flavour of a Strichartz inequality, from which we derive a decay estimate of  $L^p$ -norms for  $p \leq \frac{d^2}{d-1}$ . The proof is given in Sections 3 and 4. We explain how the strategy extends to general fluxes  $f$  in Section 5.

**Acknowledgements.** I am indebted to C. Dafermos, whose precious comments helped me to improve this article, in particular in giving full credit to previous contributors. I also thank L. Silvestre for correcting a miscalculation.

## 2 The multi-d Burgers equation

For a conservation law of the general form (1), it is harmless to assume  $f(0) = 0$ . By choosing an appropriate inertial frame, which does not affect the norms  $\|u(t)\|_p$ , we may also assume  $f'(0) = 0$ . Thus  $f(s) = O(s^2)$  at the origin. Say that  $f(s) \sim s^k \vec{v}_1$  as  $s \rightarrow 0$ , where  $\vec{v}_1$  is a non-zero constant vector. We may perform a linear change of the spacial coordinates such that  $f_1(s) \sim \frac{s^k}{k}$  and  $f_j(s) = o(|s|^k)$  otherwise. Unless we meet a flat component, the process can be continued until we find a new coordinate system  $(y_1, \dots, y_n)$  in which

$$f_j(s) \sim \frac{s^{k_j}}{k_j}, \quad 2 \leq k_1 < k_2 \cdots < k_n.$$

Generically, we have  $k_j = j + 1$  for every  $j \in \llbracket 1, n \rrbracket$ . This is the reason why we consider from now on the following scalar conservation law, which we call the *multi-dimensional Burgers equation* :

$$(6) \quad \partial_t u + \partial_j \frac{u^2}{2} + \cdots + \partial_n \frac{u^{n+1}}{n+1} = 0.$$

This particular flux was already considered by G. Crippa et al. [3]. If  $n = 1$ , we recognize the original Burgers equation. The equation (6) is a prototype for genuinely nonlinear conservation laws, those which satisfy the assumption

$$(7) \quad \det(f'', \dots, f^{(n+1)}) \neq 0.$$

The latter condition is a variant of the *non-degeneracy condition* at work in the kinetic formulation of the equation (1) ; see [8] or [9].

Let us review two preliminary answers to our natural question, in the context of (6).

- On the one hand, we might assume that  $u_0 \in L^1 \cap L^p(\mathbb{R}^n)$  for some  $p > 1$ . Let us define

$$u_{0m} := \max(-m, \min(u_0, m)) \in L^1 \cap L^\infty(\mathbb{R}^n),$$

which tends towards  $u_0$  in the  $L^1$ -norm. We have  $u = \lim_{m \rightarrow +\infty} u_m$ , where  $u_m$  is the solution associated with the data  $u_{0m}$ , and the limit stands in  $C_b(\mathbb{R}_+; L^1(\mathbb{R}^n))$ . Because of  $\|u_m(t)\|_p \leq \|u_{0m}\|_p \leq \|u_0\|_p$ , the sequence  $(u_m)_{m \geq 1}$  is bounded in  $L^\infty(\mathbb{R}_+; L^p(\mathbb{R}^n))$ . We infer that  $u \in L^\infty(\mathbb{R}_+; L^p(\mathbb{R}^n))$ , and  $u_m$  converges towards  $u$  in  $L^\infty(\mathbb{R}_+; L^q(\mathbb{R}^n))$  for every  $q \in [1, p)$ . In addition,  $u_m$  converges weakly in  $L^\infty(\mathbb{R}_+; L^p(\mathbb{R}^n))$ . If  $p > n + 1$ , we may pass to the limit as  $m \rightarrow +\infty$  in the sequences

$$u_m^k, \quad \text{sgn}(u_m - a)(u_m^k - a^k).$$

Passing to the limit in the integral formulations (2) and (3), we conclude that  $u$  is a genuine entropy solution of the Cauchy problem. Notice that the argument does not work out when  $p = n + 1$ , because of the last component of the flux: we are not certain that  $u_m^{n+1}$  converges in  $L^1_{\text{loc}}$  towards  $u^{n+1}$ . If  $p > n + 2$ , we find as well that  $u$  satisfies the entropy inequality for the pair  $(\bar{\eta}, \bar{q})$ .

The drawback of this argument is that it does not exploit the nonlinearity of the equation, a property which is expected to imply some kind of regularization or dispersion (see Theorem 4 and Proposition 1 of [8]). We should be able to lower somehow the threshold  $p > n + 1$ .

- The other answer concerns the one-dimensional case ( $n = 1$ ). The Kruzhkov solution of the classical Burgers equation satisfies the inequality

$$(8) \quad TV \left( \frac{u(t)^2}{2} \right) \leq \frac{2\|u_0\|_1}{t},$$

due to Bénilan & Crandall [1], who exploit the homogeneity of the flux. It is extended by Dafermos [4] to situations where the flux  $f$  has an inflexion point and the data  $u_0$  has bounded variations, by a careful use of the generalized backward characteristics. It implies in particular an estimate

$$(9) \quad \|u(t)\|_\infty \leq 2 \sqrt{\frac{2\|u_0\|_1}{t}}.$$

This shows that the assumption  $u_0 \in L^1(\mathbb{R})$  is sufficient in order that  $u$  be a true entropy solution. This is definitely better than the threshold  $L^1 \cap L^2(\mathbb{R}^n)$  considered in the previous paragraph.

Dafermos' argument, which is the most general one, uses the ordered structure of the real line. Backward characteristics are not unique in general. Given a base point  $(x^*, t^*)$  in the upper half-plane, one has to define and analyse the minimal and the maximal ones. These notions have not yet been extended to the multi-dimensional situation (see however [12] for a weaker notion).

Our main result here is the following statement. It tells us that  $L^1(\mathbb{R}^n)$  is the right space for initial data.

**Theorem 2.1 (Multi-d Burgers equation.)** *Suppose  $u_0 \in L^1(\mathbb{R}^n)$ . Define  $u(t) = S_t u_0$  and set  $u(t, y) = u(t)(y)$  for  $t > 0$  and  $y \in \mathbb{R}^n$ . Then*

1. *There holds an algebraic decay:*

$$(10) \quad \|u(t)\|_{\frac{d^2}{d-1}} \leq c_d t^{-\delta} \|u_0\|_1^\gamma,$$

where

$$\gamma = \frac{d^2 + 1}{d(d^2 - d + 2)}, \quad \delta = 2 \frac{(d-1)(d^2 - d + 1)}{d^2(d^2 - d + 2)} < 1.$$

2. *For every  $k \in \llbracket 1, d+1 \rrbracket$ , there holds  $u^k \in L^1_{\text{loc}}(\mathbb{R}_+ \times \mathbb{R}^n)$ .*

3. *The function  $u$  is an entropy solution of the Cauchy problem.*

4. *It satisfies the additional entropy inequality*

$$\partial_t \bar{\eta}(u) + \text{div}_y \bar{q}(u) \leq 0.$$

5. *If in addition  $u_0 \in L^1 \cap L^d(\mathbb{R}^n)$ , then  $u \in L^{\frac{d^2}{d-1}}(\mathbb{R}_+ \times \mathbb{R}^n)$  and there holds*

$$(11) \quad \left( \int_0^\infty \int_{\mathbb{R}^n} u^{\frac{d^2}{d-1}} dy dt \right)^{\frac{d-1}{d}} \leq c_d \left( \int_{\mathbb{R}^n} u_0(y)^d dy \right)^{\frac{1}{2}} \left( \int_{\mathbb{R}^n} u_0(y) dy \right)^{\frac{1}{2}}.$$

### Comments.

- The assumption that  $u_0 \in L^1(\mathbb{R}^n)$  extends that available in the 1-dimension situation. However, when  $n = 1$ , Theorem 2.1 provides an estimate of  $u(t)$  in  $L^4(\mathbb{R})$  only, instead of the known  $L^\infty(\mathbb{R})$  or  $BV(\mathbb{R})$ . Our results are new only when  $n \geq 2$ .

- The decay result is optimal when  $n = 1$ , where it states that

$$\|u(t)\|_4 \leq \text{cst} \cdot t^{-3/8} \|u_0\|_1^{5/8}.$$

This is the exact rate for an N-wave

$$N_L(t, y) = \begin{cases} \frac{y}{1+t} & \text{if } y \in (0, L\sqrt{1+t}), \\ 0 & \text{otherwise.} \end{cases}$$

It raises therefore the question whether the decay rate given by (10) is accurate also when  $n \geq 2$ .

- Estimate (11) resembles a Strichartz inequality. It seems to be new in this situation where the principal part is not a linear operator, but a quasilinear one.
- By Hölder interpolation, together with  $u \in L^\infty(\mathbb{R}_+; L^1(\mathbb{R}^n))$ , (10) implies

$$(12) \quad \|u(t)\|_q \leq c_{d,q} t^{-\kappa/q'} \|u_0\|_1^{1-\nu/q'}, \quad \forall q \in (1, \frac{d^2}{d+1}),$$

where

$$\kappa = 2 \frac{d-1}{d^2-d+2} \quad \text{and} \quad \nu = \frac{d(d-1)}{d^2-d+2}.$$

- A natural question is whether (12) extends to every  $q \in [1, \infty]$ . In particular, is it true that  $u(t) \in L^\infty(\mathbb{R}^n)$  for every  $t > 0$ ? We now that it is true when  $n = 1$ , see (9).
- A useful contribution in this direction was obtained recently by L. Silvestre [12], whose Theorem 1.5 tells in particular that if  $u_0 \in L^1 \cap L^\infty(\mathbb{R}^n)$ , then

$$\|u(t)\|_\infty \leq C(\|u_0\|_\infty; \mu) \|u_0\|_1^\mu t^{-n\mu}$$

for every  $\mu < \mu_0$  where

$$\mu_0 = \frac{2}{d^2-d+2}.$$

This decay is almost the same as that suggested by extrapolation to  $q = \infty$  of ours, because of

$$n\mu_0 = \kappa = \frac{\kappa}{q'}.$$

It would be exactly that one if the limit exponent  $\mu_0$  was allowed, and the dependency of the constant upon  $\|u_0\|_\infty$  was removed.



## 2.1 Other “monomial” scalar conservation laws

As suggested above, we may be interested into more general conservation laws, whose fluxes are monomial. Denoting  $P_k(s) = \frac{s^k}{k}$ , consider the PDE

$$(13) \quad \partial_t u + \partial_1 P_{k_1}(u) + \cdots + \partial_n P_{k_n}(u) = 0,$$

where  $1 < k_1 < \cdots < k_n$  are integers. We leave, as a tedious exercise, the interesting reader to adapt the calculations of the two next sections to (13), to prove the following result. We denote

$$K = \sum_{j=1}^n k_j, \quad N = 1 + 2K - n.$$

**Theorem 2.2** *Suppose that  $nk_n < N$ . Then for every initial data  $u_0 \in L^1(\mathbb{R}^n)$ , the abstract solution given by the continuous extension of the semi-group  $(S_t)_{t \geq 0}$  to  $L^1(\mathbb{R}^n)$ , is actually an entropy solution of the Cauchy problem for (13). It satisfies a dispersion estimate*

$$\left( \int_0^\infty dt \int_{\mathbb{R}^n} u(t, y)^{\frac{N}{n}} dy \right)^{\frac{n}{n+1}} \leq c_d \left( \int_{\mathbb{R}^n} u_0(y) dy \right)^{1-\theta} \left( \int_{\mathbb{R}^n} u_0(y)^{k_n} dy \right)^\theta,$$

where

$$\theta = \frac{K - n}{(n+1)(k_n - 1)}.$$

It decays as follows

$$\|u(t)\|_{\frac{N}{n}} \leq c_d \frac{\|u_0\|_1^\gamma}{t^\delta}$$

where

$$\gamma = \frac{n}{N} + \frac{N - n}{N(N - K)}, \quad \delta = \frac{n(N - n)}{N(N - K)}.$$

The rôle of the assumption  $nk_n < N$  is to allow us to estimate  $\|u(t)\|_{k_n}$  in terms of  $\|u(t)\|_1$  and  $\|u(t)\|_{N/n}$ , in order to apply a Gronwall argument to the dispersion estimate. Notice that it is always satisfied in one space dimension, because then  $1 \cdot k_1 < N = 2k_1$

**Remark.** If  $k_n$  is larger than  $\frac{N}{n}$ , there should be a weaker result. There will be some exponent  $p = p(k_n, N) \in (1, k_n)$  such that if  $u_0 \in L^1 \cap L^p(\mathbb{R}^n)$ , then the abstract solution is actually an entropy solution. We leave the calculation of  $p(k_n, N)$  to the motivated reader.

### 3 Proof of Estimate (11)

Because  $u$  is obtained as the limit in  $C(\mathbb{R}_+; L^1(\mathbb{R}^n))$  of  $u_m$ , the solution associated with the data  $u_{0m} = \text{Proj}_{[-m,m]} u_0$ , the estimates (10) and (11) need only to be proved when the initial data belongs to  $L^1 \cap L^\infty(\mathbb{R}^n)$ , that is within Kruzhkov's theory. Then they extend to  $L^1$ -data by a density argument.

When  $u_0 \in L^1 \cap L^\infty$ , (11) will provide a uniform bound of

$$\int_0^\infty \|u_m(t)\|_{\frac{d^2}{d-1}} dt.$$

Then, because of  $u_m \rightarrow u$  in  $C(\mathbb{R}_+; L^1(\mathbb{R}^n))$  as  $m \rightarrow +\infty$ , we infer by interpolation that the convergence holds true in every space  $L^q(\mathbb{R}_+; L^p(\mathbb{R}^n))$  for which

$$1 \leq p < \frac{d^2}{d-1} \quad (\text{equivalently } \frac{d^2}{d-1} < q \leq \infty) \quad \text{and} \quad \frac{d-1}{p} + \frac{d^2-d+1}{q} = d-1.$$

Because of  $(S_t u_0)_\pm \leq S_t(u_0)_\pm$ , it is enough to consider data that are either non-negative or non-positive. But since  $v(t, y) = -u(t, -y_1, y_2, \dots, (-1)^n y_n)$  is the entropy solution associated with  $v_0(y) = -u_0(-y_1, y_2, \dots, (-1)^n y_n)$ , it suffices to prove (11) for non-negative data and solutions. We therefore assume from now on that  $u_0 \geq 0$ , and thus  $u \geq 0$  over  $\mathbb{R}_+ \times \mathbb{R}^n$ .

If  $a \in \mathbb{R}$ , we define a symmetric matrix

$$M(a) = \left( \frac{a^{i+j+1}}{i+j+1} \right)_{0 \leq i, j, \leq n}.$$

Remarking that

$$M(a) = \int_0^a V(s) \otimes V(s) ds, \quad V(s) = \begin{pmatrix} 1 \\ \vdots \\ s^n \end{pmatrix},$$

we obtain that  $M(a)$  is positive definite whenever  $a > 0$ . We have obviously

$$\det M(a) = H_d a^{d^2},$$

where

$$H_d = \left\| \frac{1}{i+j+1} \right\|_{0 \leq i, j, \leq n} > 0$$

is the determinant of the Hilbert matrix (this is the only case where we do not write  $c_d$  for a dimensional constant).

Let us form the symmetric tensor

$$T(t, y) = M(u(t, y)),$$

with positive semi-definite values. Its first line is formed of  $(u, f(u))$  and therefore is divergence-free by (6). The second line is formed of  $(\bar{\eta}(u), \bar{q}(u))$ , an entropy-flux pair. It is not divergence-free in general, although it is so away from shock waves and other singularities of the solution  $u$ . But the entropy inequality tells us that the opposite of its divergence is a non-negative, hence bounded measure,

$$\mu_1 = -\operatorname{div}_{t,y}(\bar{\eta}(u), \bar{q}(u)) \geq 0.$$

The total mass of  $\mu_1$  over a slab  $(0, \tau) \times \mathbb{R}^n$  is given by

$$\|\mu_1\| = \int_{\mathbb{R}^n} \bar{\eta}(u_0(y)) dy - \int_{\mathbb{R}^n} \bar{\eta}(u(\tau, y)) dy \leq \int_{\mathbb{R}^n} \bar{\eta}(u_0(y)) dy.$$

Notice that the latter bound does not depend of  $\tau$ . The same situation occurs for the other lines of  $T$ . They are of the form  $(\eta(u), q(u))$  where  $(\eta, q)$  is an entropy-flux pair with  $\eta$  convex over  $\mathbb{R}_+$  (recall that  $u$  takes only non-negative values). The distribution

$$\mu = -\operatorname{div}_{t,y}(\eta(u), q(u))$$

is therefore again a bounded measure, whose total mass over  $\mathbb{R}_+ \times \mathbb{R}^n$  is bounded by

$$\int_{\mathbb{R}^n} \eta(u_0(y)) dy.$$

We conclude that the row-wise divergence of  $T$  is a (vector-valued) bounded measure, whose total mass is bounded above by

$$\sum_{j=2}^d \int_{\mathbb{R}^n} \frac{u_0(y)^j}{j} dy.$$

We may therefore apply Compensated integrability (Theorems 2.2 and 2.3 of [11]) to the tensor  $T$ , that is

$$\int_0^\tau dt \int_{\mathbb{R}^n} (\det T)^{\frac{1}{d-1}} dy \leq c_d \left( \|T_{0\bullet}(0, \cdot)\|_1 + \|T_{0\bullet}(\tau, \cdot)\|_1 + \|\operatorname{Div}_{t,y} T\|_{\mathcal{M}((0,\tau) \times \mathbb{R}^n)} \right)^{\frac{d}{d-1}}.$$

Because of

$$\|T_{0\bullet}(t, \cdot)\|_1 \leq \sum_{j=1}^d \int_{\mathbb{R}^n} \frac{u(t, y)^j}{j} dy \leq \sum_{j=1}^d \int_{\mathbb{R}^n} \frac{u_0(y)^j}{j} dy,$$

we deduce

$$(14) \quad \int_0^\tau dt \int_{\mathbb{R}^n} u^{\frac{d^2}{d-1}} dy \leq c_d \left( \sum_{j=1}^d \int_{\mathbb{R}^n} u_0(y)^j dy \right)^{\frac{d}{d-1}}.$$

The only bad feature in the estimate (14) is the lack of homogeneity of its right-hand side. To recover a well-balanced inequality, we exploit an idea already used in [10]. We begin by remarking that if  $\lambda > 0$  is a constant parameter, then the function

$$v(t, y) = \frac{1}{\lambda} u(\lambda t, \lambda^2 y_1, \dots, \lambda^d y_n)$$

is the entropy solution associated with the initial data

$$v_0(y) = \frac{1}{\lambda} u_0(\lambda^2 y_1, \dots, \lambda^d y_n).$$

Applying (14) to the pair  $(v, v_0)$  instead, then using

$$\int_{\mathbb{R}^n} v_0(y)^j dy = \lambda^{-j - \frac{(d-1)(d+2)}{2}} \int_{\mathbb{R}^n} u_0(y)^j dy$$

and

$$\left( \int_0^\infty \int_{\mathbb{R}^n} v^{\frac{d^2}{d-1}} dy dt \right)^{\frac{d-1}{d}} = \lambda^{-d - \frac{d^2-1}{2}} \left( \int_0^\infty \int_{\mathbb{R}^n} u^{\frac{d^2}{d-1}} dy dt \right)^{\frac{d-1}{d}},$$

we get a parametrized inequality

$$\left( \int_0^\tau dt \int_{\mathbb{R}^n} u^{\frac{d^2}{d-1}} dy \right)^{\frac{d-1}{d}} \leq c_d \lambda^{\frac{d+1}{2}} \sum_{j=1}^d \lambda^{-j} \int_{\mathbb{R}^n} u_0(y)^j dy.$$

In order to minimize the right-hand side, we choose the value

$$\lambda = \left( \int_{\mathbb{R}^n} u_0(y)^d dy / \int_{\mathbb{R}^n} u_0(y) dy \right)^{\frac{1}{d-1}}.$$

The extreme terms, for  $j = 1$  or  $d$ , contribute on an equal foot with

$$\left( \int_{\mathbb{R}^n} u_0(y)^d dy \right)^{\frac{1}{2}} \left( \int_{\mathbb{R}^n} u_0(y)^2 dy \right)^{\frac{1}{2}}.$$

The other ones, which are

$$\left( \int_{\mathbb{R}^n} u_0(y)^d dy / \int_{\mathbb{R}^n} u_0(y) dy \right)^{\frac{d+1-2j}{2(d-1)}} \int_{\mathbb{R}^n} u_0^j dy,$$

are bounded by the same quantity, because of Hölder inequality. We end therefore with the fundamental estimate (11).

## 4 Proof of Theorem 2.1

We now complete the proof of our main theorem.

### 4.1 The decay result

We keep working with the assumptions  $u_0 \in L^1 \cap L^\infty(\mathbb{R}^n)$  and  $u_0 \geq 0$ .

Let us define

$$X(t) := \int_{\mathbb{R}^n} u^{\frac{d^2}{d-1}} dy = \|u(t)\|_{\frac{d^2}{d-1}}^{\frac{d^2}{d-1}}.$$

From the Hölder inequality, we have

$$\|u_0\|_d \leq \|u_0\|_1^{1-\theta} \|u_0\|_{\frac{d^2}{d-1}}^\theta, \quad \theta = \frac{d(d-1)}{d^2-d+1} \in (0, 1).$$

The inequality (11) implies therefore

$$\int_0^\infty X(t) dt \leq c_d \|u_0\|_1^\alpha X(0)^\beta$$

where

$$\beta = \frac{\theta}{2}, \quad \alpha = \frac{d(d^2+1)}{2(d-1)(d^2-d+1)}.$$

Considering the solution  $v(t, y) = u(t + \tau, y)$ , whose initial data is  $u(\tau, \cdot)$ , we also have

$$(15) \quad \int_\tau^\infty X(t) dt \leq c_d \|u(\tau)\|_1^\alpha X(\tau)^\beta \leq c_d \|u_0\|_1^\alpha X(\tau)^\beta.$$

Let us denote

$$Y(\tau) := \int_\tau^\infty X(t) dt.$$

We recast (15) as

$$Y^{\frac{1}{\beta}} + c_d \|u_0\|_1^{\frac{\alpha}{\beta}} Y' \leq 0.$$

Multiplying by  $Y^{-1/\beta}$  and integrating, we infer (mind that  $1 - \frac{1}{\beta}$  is negative)

$$t + c_d \|u_0\|_1^{\frac{\alpha}{\beta}} Y(0)^{1-\frac{1}{\beta}} \leq c_d \|u_0\|_1^{\frac{\alpha}{\beta}} Y(t)^{1-\frac{1}{\beta}}.$$

This provides a first decay estimate

$$Y(t) \leq c_d \|u_0\|_1^{\frac{\alpha}{1-\beta}} t^{-\frac{\beta}{1-\beta}}.$$

Remarking that  $t \mapsto X(t)$  is a non-increasing function, so that

$$\frac{\tau}{2} X(\tau) \leq Y\left(\frac{\tau}{2}\right),$$

we deduce the ultimate decay result

$$X(t) \leq c_d \|u_0\|_1^{\frac{\alpha}{1-\beta}} t^{-\frac{1}{1-\beta}}.$$

Restated in terms of a Lebesgue norm of  $u(t)$ , it says

$$(16) \quad \|u(t)\|_{\frac{d^2}{d-1}} \leq c_d \|u_0\|_1^\gamma t^{-\delta},$$

where

$$\gamma = \frac{d^2 + 1}{d(d^2 - d + 2)}, \quad \delta = 2 \frac{(d-1)(d^2 - d + 1)}{d^2(d^2 - d + 2)}.$$

## 4.2 The function $u$ is an entropy solution

We already know that the functions  $u_m$  are entropy solutions, with initial data  $u_{0m} \in L^1 \cap L^\infty(\mathbb{R}^n)$ . Because of (11), we have seen that  $u_m$  converges towards  $u$  in the norm of  $L^q(\mathbb{R}_+; L^p(\mathbb{R}^n))$  whenever

$$1 \leq p < \frac{d^2}{d-1} \quad \text{and} \quad \frac{d-1}{p} + \frac{d^2-d+1}{q} = d-1.$$

Because of  $\frac{d^2}{d-1} > d+1$ , this implies that  $(u_m)^k \rightarrow u^k$  for every  $k \in [[1, d+1]]$ . Recalling that

$$f(s) = \left(s, \frac{s^2}{2}, \dots, \frac{s^d}{d}\right) \quad \text{and} \quad \bar{q}(s) = \left(\frac{s^3}{3}, \dots, \frac{s^{d+1}}{d+1}\right),$$

we see that we may pass to the limit as  $m \rightarrow +\infty$  in the identity

$$\int_0^\infty dt \int_{\mathbb{R}^n} (u_m \partial_t \phi + f(u_m) \cdot \nabla_y \phi) dy + \int_{\mathbb{R}^n} u_{0m}(y) \phi(0, y) dy = 0, \quad \forall \phi \in \mathcal{D}(\mathbb{R}^{1+n}),$$

as well as in the Kruzhkov inequalities

$$\int_0^\infty dt \int_{\mathbb{R}^n} (\eta_a(u_m) \partial_t \phi + q_a(u_m) \cdot \nabla_y \phi) dy + \int_{\mathbb{R}^n} \eta_a(u_{0m}(y)) \phi(0, y) dy \geq 0, \quad \forall \phi \in \mathcal{D}^+(\mathbb{R}^{1+n}),$$

and in the inequality

$$\int_0^\infty dt \int_{\mathbb{R}^n} (\bar{\eta}(u_m) \partial_t \phi + \bar{q}(u_m) \cdot \nabla_y \phi) dy + \int_{\mathbb{R}^n} \bar{\eta}(u_{0m}(y)) \phi(0, y) dy \geq 0, \quad \forall \phi \in \mathcal{D}^+(\mathbb{R}^{1+n}).$$

Therefore  $u$  is an entropy solution with initial data  $u_0$ , which satisfies in addition the entropy inequality for the pair  $(\bar{\eta}, \bar{q})$ .

**Remark.** When  $n \geq 2$ , the Compensated Integrability cannot be applied directly to the solution  $u$ , when the data is only integrable. Because we don't know whether the  $j$ th line of  $T$  is locally integrable if  $j = 3, \dots, n+1$ ; its last component is  $\frac{u^{n+j}}{n+j}$ , where the exponent  $n+j$  is larger than  $\frac{d^2}{d-1}$ .

## 5 The strategy for general fluxes $f$

We come back to the study of a multi-dimensional conservation law of general form (1). Following the ideas developed in the Burgers case, we begin by considering a signed, bounded initial data:  $u_0 \in L^1 \cap L^\infty(\mathbb{R}^n)$ ,  $u_0 \geq 0$ . If  $a \in \mathbb{R}_+$ , we define a symmetric matrix

$$M(a) = \int_0^a F'(s) \otimes F'(s) ds,$$

where  $F(s) = (s, f_1(s), \dots, f_n(s))$ . This matrix is positive definite under the non-degeneracy condition that  $F([0, a])$  is not contained in an affine hyperplane. We denote

$$\Delta(a) := (\det M(a))^{\frac{1}{n}} \geq 0.$$

Let us define  $T(t, y) := M(u(t, y))$ . Because  $u \in L^\infty(0, \tau; L^1 \cap L^\infty(\mathbb{R}^n))$ , the tensor  $T$  is integrable over  $(0, \tau) \times \mathbb{R}^n$ . The first line of  $T$  is divergence-free. The other lines are made of

entropy-entropy flux pairs  $(f_i, Q_i)$ . Since  $f_i$  might not be convex, we cannot estimate the measure  $\mu_i = -\partial_t f_i(u) - \operatorname{div}_y Q_i(u)$  directly by the integral of  $f_i(u_0)$ . To overcome this difficulty, we define a convex function  $\phi$  over  $\mathbb{R}_+$  by

$$\phi(0) = \phi'(0) = 0, \quad \phi''(s) = |f''(s)|.$$

Remark that  $|f'| \leq \phi'$  and  $|f| \leq \phi$ . Let  $\Phi$  be the entropy flux associated with the entropy  $\phi$ . Then the measure  $\mathbf{v} := -\partial_t \phi(u) - \operatorname{div}_y \Phi(u)$  is non-negative and a bound of its total mass is as usual

$$\|\mathbf{v}\| \leq \int_{\mathbb{R}^n} \phi(u_0(y)) dy.$$

We now use the kinetic formulation of (1), a notion for which we refer to [9], Theorem 3.2.1. Recall the definition of the kinetic function  $\chi(\xi; a)$ , whose value is  $\operatorname{sgn} a$  if  $\xi$  lies between 0 and  $a$ , and is 0 otherwise. There exists a non-negative bounded measure  $m(t, y, \xi)$  such that the function  $g(t, y, \xi) = \chi(\xi; u(t, y))$  satisfies

$$\partial_t g + f'(\xi) \cdot \nabla_y g = \frac{\partial}{\partial \xi} m, \quad g(0, y; \xi) = \chi(\xi; u_0(y)).$$

If  $(\eta, q)$  is an entropy-entropy flux pair, then the measure  $\mu = -\partial_t \eta - \operatorname{div}_y q$  is given by

$$\mu = \int_{\mathbb{R}} \eta''(\xi) dm(\xi).$$

We deduce that the vector-valued measure  $\mu = (\mu_1, \dots, \mu_n)$  satisfies  $|\mu| \leq \mathbf{v}$ . This yields the estimate

$$\|\mu\| \leq \int_{\mathbb{R}^n} \phi(u_0(y)) dy.$$

We may therefore apply the compensated integrability, which gives here

$$\int_0^\tau dt \int_{\mathbb{R}^n} \Delta(u(t, y)) dy \leq c_d \left( \|F(u_0)\|_1 + \|F(u(\tau))\|_1 + \int_{\mathbb{R}^n} \phi(u_0(y)) dy \right)^{1+\frac{1}{n}}.$$

Because of  $|f| \leq \phi$  and  $\|\phi(u(\tau))\|_1 \leq \|\phi(u_0)\|_1$ , we end up with an analog of (11)

$$(17) \quad \int_0^\infty dt \int_{\mathbb{R}^n} \Delta(u(t, y)) dy \leq c_d (\|u_0\|_1 + \|\phi(u_0)\|_1)^{1+\frac{1}{n}}.$$

To improve the inequality above, we use again a scaling argument. However, because the components  $f_j$  of the flux are not homogeneous anymore, we modify simultaneously the solution and the flux, using the fact that the constant  $c_d$  in (17) does not depend upon  $f$ . Our new



dependent variables are

$$v(t, y) = \frac{1}{\lambda} u(\lambda t, Py), \quad v_0(y) = v(0, y) = \frac{1}{\lambda} u_0(Py)$$

where  $P \in \mathbf{GL}_n(\mathbb{R})$  is a matrix to be chosen later. The function  $v$  is an entropy solution of the Cauchy problem associated with the conservation law

$$\partial_t v + \operatorname{div}_y g(v) = 0$$

for the flux

$$(18) \quad g(s) = P^{-1} f(\lambda s).$$

The symmetric matrix  $N(a)$  that plays the role of  $M(a)$  for (18) is given by the formula

$$N(s) = \lambda R M(\lambda s) R^T, \quad R = \operatorname{diag}\left(\frac{1}{\lambda}, P^{-1}\right).$$

We have  $\det N(s) = \lambda^{n-1} (\det P)^{-2} \det M(\lambda s)$ , from which we derive

$$\int_0^\infty \int_{\mathbb{R}^n} (\det N(v))^\frac{1}{n} dy dt = \lambda^{-\frac{1}{n}} (\det P)^{-1-\frac{2}{n}} \int_0^\infty \int_{\mathbb{R}^n} \Delta(u) dy dt.$$

When applying (18) to  $v$  and  $g$ , one integral in the right-hand side transforms easily:

$$\int_{\mathbb{R}^n} v_0(y) dy = \frac{1}{\lambda \det P} \int_{\mathbb{R}^n} u_0(y) dy.$$

The other one involves a modified function

$$\phi_{\lambda, P}(s) = \Psi_P(\lambda s), \quad \Psi_P''(s) = |P^{-1} f''(s)|.$$

As usual  $\Psi_P$  is fixed by  $\Psi_P(0) = \Psi_P'(0) = 0$ . We have therefore

$$\int_{\mathbb{R}^n} \phi_P(v_0(y)) dy = \frac{1}{\det P} \int_{\mathbb{R}^n} \Psi_P(u_0(y)) dy.$$

All the identities above, together with (18) applied to  $(v, g)$ , yield our parametrized estimate

$$\int_0^\infty \int_{\mathbb{R}^n} \Delta(u) dy dt \leq c_d \lambda^\frac{1}{n} (\det P)^\frac{1}{n} \left( \frac{1}{\lambda} \int_{\mathbb{R}^n} u_0 dy + \int_{\mathbb{R}^n} \Psi_P(u_0) dy \right)^{1+\frac{1}{n}}.$$

We optimize this inequality with respect to  $\lambda$ , by choosing

$$\lambda = \frac{\int_{\mathbb{R}^n} u_0 dy}{\int_{\mathbb{R}^n} \Psi_P(u_0) dy}.$$

We infer

$$\int_0^\infty \int_{\mathbb{R}^n} \Delta(u) dy dt \leq c_d (\det P)^{\frac{1}{n}} \left( \int_{\mathbb{R}^n} u_0 dy \right)^{\frac{1}{n}} \int_{\mathbb{R}^n} \Psi_P(u_0) dy.$$

There remains to minimize the right-hand side with respect to  $P$  :

$$(19) \quad \int_0^\infty \int_{\mathbb{R}^n} \Delta(u) dy dt \leq c_d \left( \int_{\mathbb{R}^n} u_0 dy \right)^{\frac{1}{n}} I[u_0],$$

where

$$I[w] := \inf_{P \in \mathbf{GL}_n(\mathbb{R})} (\det P)^{\frac{1}{n}} \int_{\mathbb{R}^n} \Psi_P(w(y)) dy.$$

The calculation of  $I[w]$  has to be made on a case-by-case basis.

Let us define again

$$X(t) = \int_{\mathbb{R}^n} \Delta(u(t, y)) dy.$$

Applying (19) on an interval  $(\tau, \infty)$  instead, and using the decay of the  $L^1$ -norm, we arrive to

$$\int_\tau^\infty X(t) dt \leq c_d \left( \int_{\mathbb{R}^n} u_0 dy \right)^{\frac{1}{n}} I[u(\tau)].$$

A decay result will be obtained through a Gronwall argument, whenever we can estimate  $I[u(t)]$  in terms of  $\|u(t)\|_1$  and  $X(t)$ .

## References

- [1] Ph. Bénylan, M. G. Crandall. Regularizing effects of homogeneous evolution equations. *Contributions to analysis and geometry (Baltimore, Md., 1980)*. Johns Hopkins Univ. Press, Baltimore, Md. (1981), pp 23–39.
- [2] M. Crandall. The semigroup approach to first order quasilinear equations in several space variables. *Israel J. Math.*, **12** (1972), pp 108–132.

- [3] G. Crippa, F. Otto, M. Westdickenberg. Regularizing effect of nonlinearity in multidimensional scalar conservation laws. *Transport equations and multi-D hyperbolic conservation laws*, Lect. Notes Unione Mat. Ital., **5**, Springer, Berlin, (2008), pp 77–128.
- [4] C. Dafermos. Regularity and large time behaviour of solutions of a conservation law without convexity. *Proc. Royal Soc. Edinburgh*, **99A** (1985), pp 201–239.
- [5] F. Golse. Nonlinear regularizing effect for conservation laws. In *Hyperbolic problems: theory, numerics and applications*. Proc. Sympos. Appl. Math. 67, Part 1, Amer. Math. Soc., Providence, RI, (2009), pp 73–92.
- [6] F. Golse, B. Perthame. Optimal regularizing effect for scalar conservation laws. *Rev. Mat. Iberoam.*, **29** (2013), pp 1477–1504.
- [7] S. Kružkov. First order quasilinear equations with several independent variables (in Russian). *Mat. Sbornik (N.S.)*, **81 (123)** (1970), pp 228–255.
- [8] P.-L. Lions, B. Perthame, E. Tadmor. A kinetic formulation of multidimensional scalar conservation laws and related equations. *J. Amer. Math. Soc.*, **7** (1994), pp 169–191.
- [9] B. Perthame. *Kinetic formulation of conservation laws*, Oxford lecture series in Math. & its Appl. **21**. Oxford (2002).
- [10] D. Serre. Divergence-free positive symmetric tensors and fluid dynamics. *Annales de l'Institut Henri Poincaré (analyse non linéaire)*. **35** (2018), pp 1209–1234. <https://doi.org/10.1016/j.anihpc.2017.11.002>.
- [11] D. Serre. Compensated integrability. Applications to the Vlasov–Poisson equation and other models in mathematical physics. *Journal de Mathématiques Pures et Appliquées*. To appear.
- [12] L. Silvestre. Oscillation properties of scalar conservation laws. Preprint arXiv:1708.03401v3.
- [13] L. Tartar. Compensated compactness and applications to partial differential equations. *Nonlinear analysis and mechanics: Heriot-Watt Symposium*, Vol. IV, Res. Notes in Math., 39, Pitman (1979), pp 136–212.